

# Perturbed Cross-flow Boundary Layer: nontrivial effects of the obliquity angle

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# Introduction

- The transient as well the long time behaviors of arbitrary **three-dimensional disturbances** acting on a cross-flow boundary layer are investigated.
- The **cross flow boundary layer** is one of the most important boundary layer in the engineering application (aerospace, mechanical, wind,..)

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# Introduction

- The transient as well the long time behaviors of arbitrary three-dimensional disturbances acting on a cross-flow boundary layer are investigated.
- The cross flow boundary layer is one of the most important boundary layer in the engineering application (aerospace, mechanical, wind,..)
- The role of the direction of the perturbation in respect to the base flow has been poorly investigated (Breuer & Kuraishi (Phys. Fluids 1994), Taylor & Peake (jfm 1998))

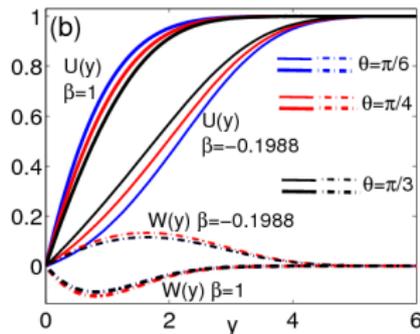
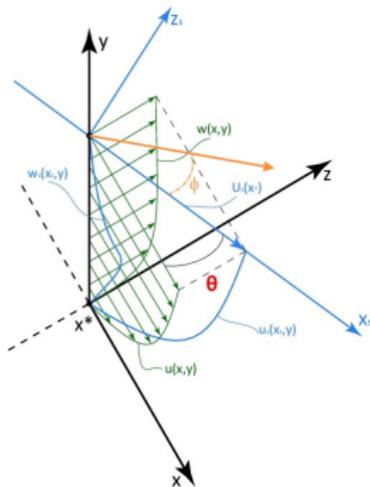
We present an exploratory analysis of the perturbation life, where a major focus is put on the obliquity of the perturbation in respect to the base flow

the stability property  
wall perturbed pressure



## Base Flows

Falkner-Skan-Cooke (FSC) velocity profiles (*Cooke, J. C. 1950 Proc. Camb. Phil. Soc.*)



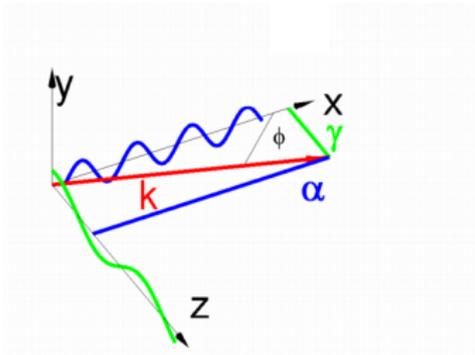
Base flow parameters:

- $Re$  Reynolds number (based on  $\delta^*$ )
- $\beta$  Hartree parameter (pressure gradient)
- $\theta$  Cross-Flow angle (between the the streamwise direction and the chordwise direction)



## Formulation

- Linearized Navier Stokes equation
- Velocity and vorticity formulation (*Criminale & Drazin, Stud. Appl. Math., 1990*);
- Laplace-Fourier transform in  $x$  and  $z$  directions:
  - $\alpha$  streamwise wavenumber
  - $\gamma$  spanwise wavenumber
  - $k$  polar wavenumber
  - $\phi$  obliquity angle between the perturbation and the streamwise direction
  - $v(y, t = 0) = y^2 \exp(-y^2), \quad \omega_y(y, t = 0) = 0$



# Measure of the transient

- Kinetic energy density  $e$ :

$$e(t; \alpha, \gamma) = \frac{1}{2} \int_{-y_d}^{+y_d} (|\hat{u}|^2 + |\hat{v}|^2 + |\hat{w}|^2) dy$$

- Amplification factor  $G$ :

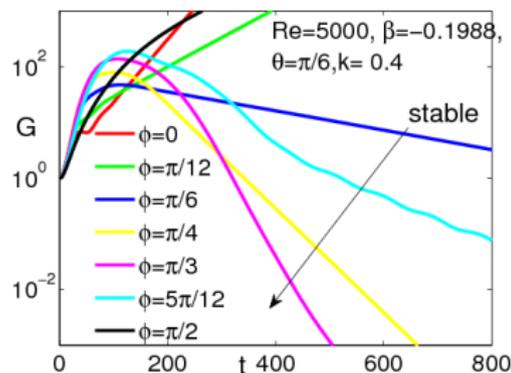
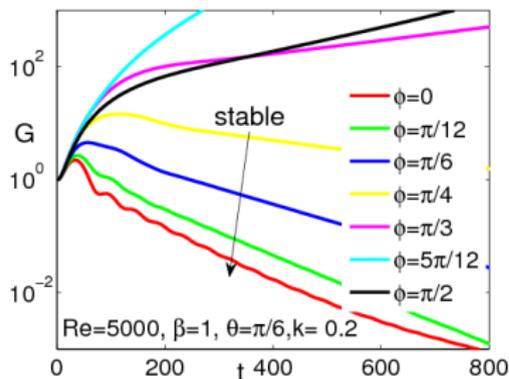
$$G(t; \alpha, \gamma) = \frac{e(t; \alpha, \gamma)}{e(t=0; \alpha, \gamma)}$$

- Temporal growth rate  $r$ :

$$r(t; \alpha, \gamma) = \frac{\log(e)}{2t}$$



# Transient behavior



- It is difficult to come across a general trend
- **Counter-intuitive behavior:** there are configuration in which longitudinal and orthogonal waves are unstable while oblique waves are stable



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# Asymptotic Growth Rate

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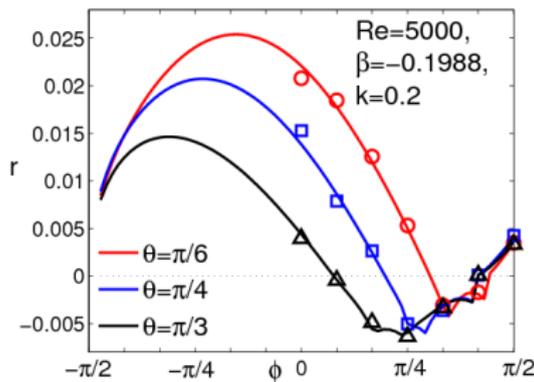
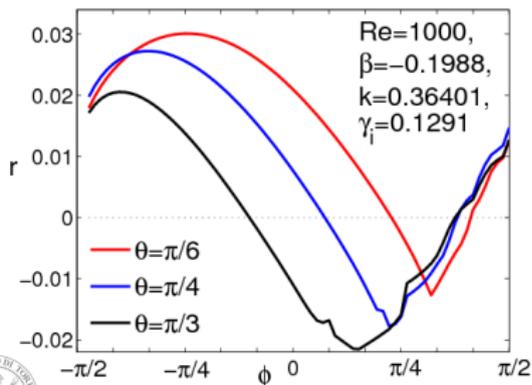
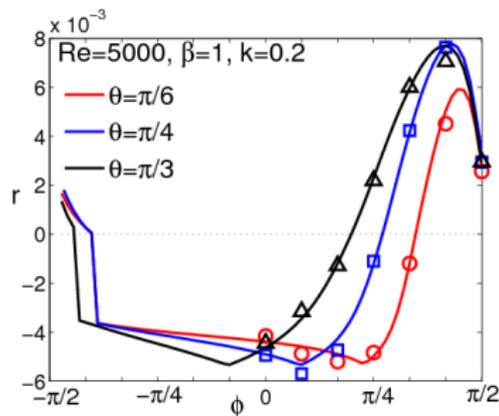
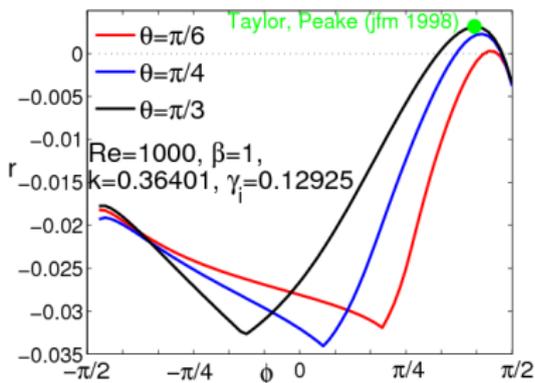
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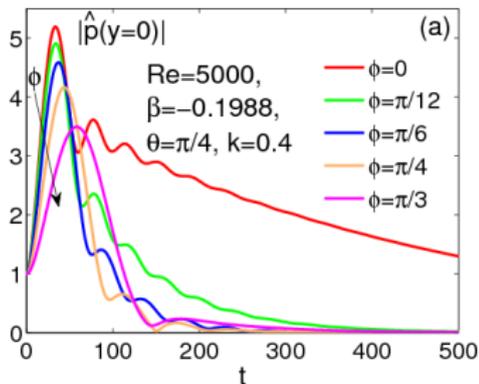
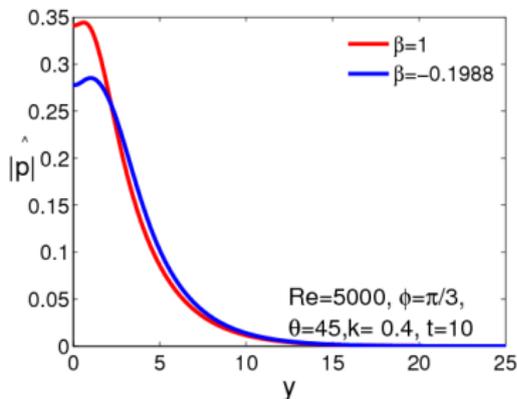


## Pressure

The pressure is computed a posteriori by the Poisson equation.

A **Pressure Amplification** is defined as

$$P = |\hat{p}(y = 0, t)| / |\hat{p}(y = 0, t = 0)|$$



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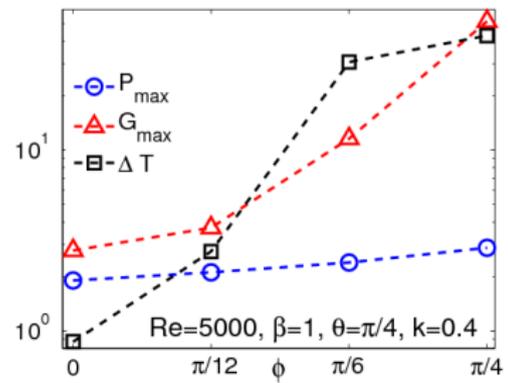
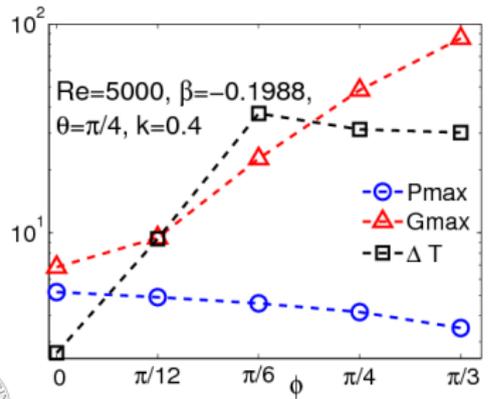
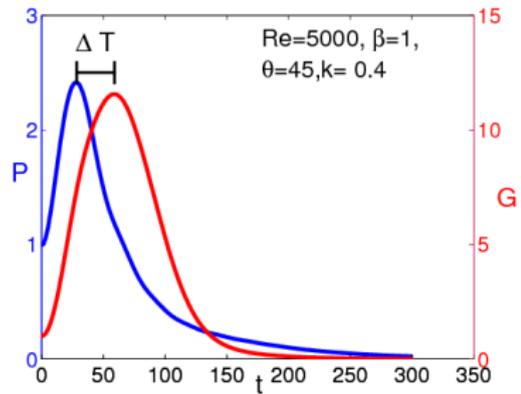
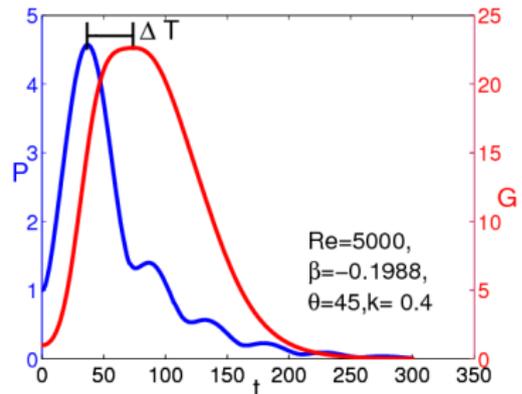
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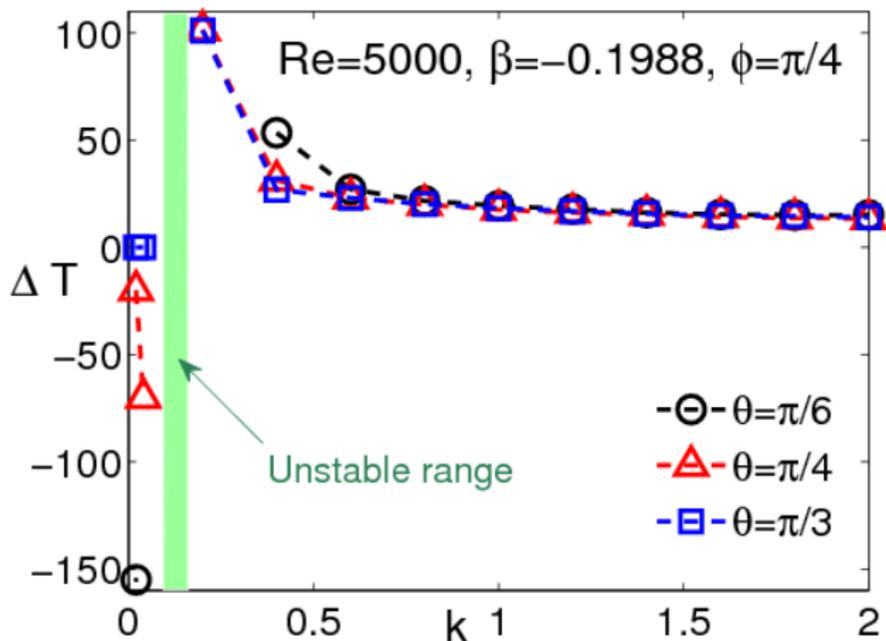
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# Pressure Transient



# Wall Pressure Delay

$$\Delta T = t(G = G_{max}) - t(P = P_{max})$$



# Conclusion

## Effects of the obliquity angle

- Let's define  $\phi_{min}$  and  $\phi_{max}$  as the obliquity angle for which the growth rate reaches its minimum/maximum value. They both decrease as  $\theta$  and  $k$  increase
- With  $\beta = 1$  the most unstable waves have about  $\phi = 5\pi/12$
- With  $\beta = -0.1988$  the most unstable waves have negative obliquity
- The perturbed pressure measured at the wall shows to be anticipated (short waves) or retarded (long waves) with respect to the kinetic energy evolution. These times increase with the obliquity angle

