

Does the Kolmogorov scaling bridge hydrodynamic linear stability and turbulence

S. Scarsoglio, F. De Santi, D. Tordella

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Does the Kolmogorov scaling bridge hydrodynamic linear stability and turbulence?

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Energy spectrum in fully developed turbulence

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- Phenomenology of turbulence Kolmogorov 1941:
 - -5/3 power-law for the energy spectrum over the inertial range;

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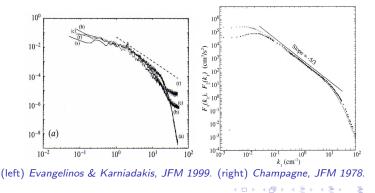
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- Phenomenology of turbulence Kolmogorov 1941: -5/3 power-law for the energy spectrum over the inertial range;
- Common criterium for the production of a fully developed turbulent field to verify such a scaling (e.g. *Frisch*, 1995; *Sreenivasan* & Antonia, ARFM, 1997; Kraichnan, Phys. Fluids, 1967).





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 - To understand how spectral representation can effectively highlight the nonlinear interaction among different scales;

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- Different typical perturbed shear systems: plane Poiseuille flow and bluff-body wake.



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- The set of small 3D perturbations:
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 - Leaves aside the nonlinear interaction among the different scales;



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• The perturbative evolution is ruled by the **initial-value problem** associated to the Navier-Stokes linearized formulation.



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• The linear transient dynamics offers a great variety of very different behaviors (*Scarsoglio et al., 2009, 2010, 2011*):



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 - \Rightarrow Understand how the energy spectrum behaves;



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 ⇒ Understand how the energy spectrum behaves;
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 ⇒ Understand how the energy spectrum behaves;
- Is the linearized perturbative system able to show a powerlaw scaling for the energy spectrum in an analogous way to the Kolmogorov argument?
- We determine the energy decay exponent of arbitrary perturbations in their asymptotic states and we compare it with the -5/3 Kolmogorov decay.



Perturbation scheme

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• Linear 3D perturbative equations in terms of velocity and vorticity (*Criminale & Drazin, Stud. Appl. Math., 1990*);

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- Linear 3D perturbative equations in terms of velocity and vorticity (*Criminale & Drazin, Stud. Appl. Math., 1990*);
- Laplace-Fourier (wake) and Fourier-Fourier (channel) transform in the x and z directions.



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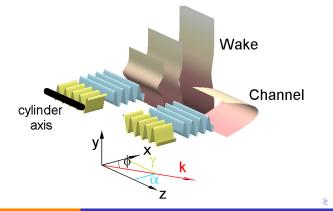
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• Perturbative linearized system: $\frac{\partial^2 \hat{v}}{\partial y^2} - k^2 \hat{v} = \hat{\Gamma}$ $\frac{\partial \hat{\Gamma}}{\partial t} = i\alpha \left(\frac{d^2 U}{dy^2} \hat{v} - U\hat{\Gamma}\right) + \frac{1}{Re} \left(\frac{\partial^2 \hat{\Gamma}}{\partial y^2} - k^2 \hat{\Gamma}\right)$ $\frac{\partial \hat{\omega}_y}{\partial t} = -i\alpha U \hat{\omega}_y - i\gamma \frac{dU}{dy} \hat{v} + \frac{1}{Re} \left(\frac{\partial^2 \hat{\omega}_y}{\partial y^2} - k^2 \hat{\omega}_y\right)$

The transversal velocity and vorticity components are \hat{v} and $\hat{\omega}_y$ respectively, $\hat{\Gamma}$ is defined as $\widetilde{\Gamma} = \partial_x \widetilde{\omega}_z - \partial_z \widetilde{\omega}_x$.



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• Initial conditions: symmetric and asymmetric inputs;



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• Perturbative linearized system:

$$\begin{array}{lll} \frac{\partial^2 \hat{v}}{\partial y^2} & - & k^2 \hat{v} = \hat{\Gamma} \\ \\ \frac{\partial \hat{\Gamma}}{\partial t} & = & i\alpha \left(\frac{d^2 U}{dy^2} \hat{v} - U \hat{\Gamma} \right) + \frac{1}{Re} \left(\frac{\partial^2 \hat{\Gamma}}{\partial y^2} - k^2 \hat{\Gamma} \right) \\ \\ \frac{\partial \hat{\omega}_y}{\partial t} & = & -i\alpha U \hat{\omega}_y - i\gamma \frac{dU}{dy} \hat{v} + \frac{1}{Re} \left(\frac{\partial^2 \hat{\omega}_y}{\partial y^2} - k^2 \hat{\omega}_y \right) \end{array}$$

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- Initial conditions: symmetric and asymmetric inputs;
- Boundary conditions: $(\hat{u}, \hat{v}, \hat{w}) \to 0$ as $y \to \pm \infty$ and at walls.



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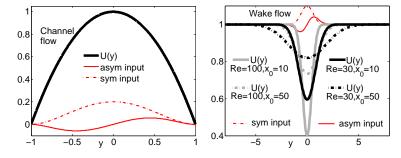
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• Kinetic energy density *e*:

$$e(t;\alpha,\gamma) = \frac{1}{2} \int_{-y_d}^{+y_d} (|\hat{u}|^2 + |\hat{v}|^2 + |\hat{w}|^2) dy$$

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• Amplification factor G:

$$G(t;\alpha,\gamma) = \frac{e(t;\alpha,\gamma)}{e(t=0;\alpha,\gamma)}$$

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• Amplification factor G:

$$G(t; \alpha, \gamma) = \frac{e(t; \alpha, \gamma)}{e(t = 0; \alpha, \gamma)}$$

• Temporal growth rate *r*:

$$r(t;\alpha,\gamma) = \frac{log[e(t;\alpha,\gamma)]}{2t}$$

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• Temporal growth rate r:

$$r(t;\alpha,\gamma) = \frac{log[e(t;\alpha,\gamma)]}{2t}$$

• Angular frequency (pulsation) ω :

$$\omega(t;y=y_0,\alpha,\gamma)=\frac{d\varphi(t;y=y_0,\alpha,\gamma)}{dt},\qquad \varphi \ \ \text{time phase}$$



Relevant transient behaviors

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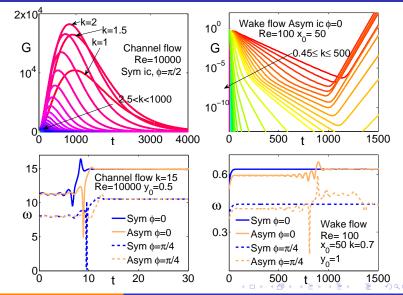
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• The energy spectrum is evaluated as the wavenumber distribution of the amplification factor, G(k);

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- The energy spectrum is evaluated as the wavenumber distribution of the amplification factor, G(k);
- The spectral representation is determined by comparing the energy of the waves when they are exiting their transient state:



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- The energy spectrum is evaluated as the wavenumber distribution of the amplification factor, G(k);
- The spectral representation is determined by comparing the energy of the waves when they are exiting their transient state;
- Every perturbation has a characteristic transient exit time, T_e;



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- The energy spectrum is evaluated as the wavenumber distribution of the amplification factor, G(k);
- The spectral representation is determined by comparing the energy of the waves when they are exiting their transient state;
- Every perturbation has a characteristic transient exit time, T_e ;
- The asymptotic condition is reached when the perturbative wave exceeds the transient exit time, T_e , that is when $r \sim const$ is satisfied for stable and unstable waves.

Scarsoglio, De Santi & Tordella, submitted to Phys. Rev. Lett., 2011.



Energy G(k) at the asymptotic state ($r \sim \text{const}$)

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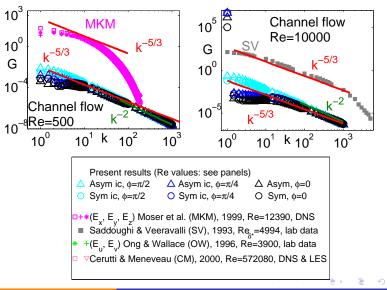
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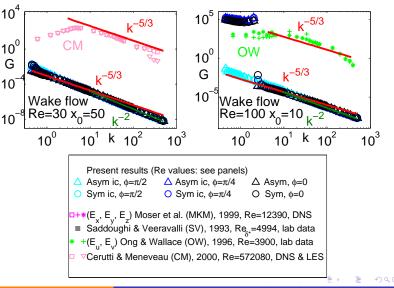
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Transient exit time $T_e(k)$

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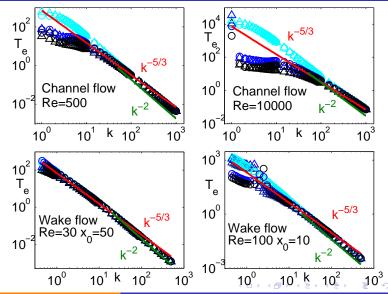
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• Spectrum determined by evaluating the energy of the waves when they are exiting their transient state;

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- Spectrum determined by evaluating the energy of the waves when they are exiting their transient state;
- Regardless the symmetry and obliquity of perturbations, there exists an intermediate range of wavenumbers in the spectrum where the energy decays with the same exponent observed for fully developed turbulent flows (-5/3), where the nonlinear interaction is considered dominant:



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- Regardless the symmetry and obliquity of perturbations, there exists an intermediate range of wavenumbers in the spectrum where the energy decays with the same exponent observed for fully developed turbulent flows (-5/3), where the nonlinear interaction is considered dominant;
- Scale-invariance of G and T_e at different (stable and unstable) Reynolds numbers and for different shear flows;
- The -5/3 spectral power-law scaling of inertial waves seems to be a general and intrinsic dynamical property of the NS solutions encompassing the nonlinear interaction.



Coming next...

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- Analysis of the perturbation transient dynamics in the 2D and 3D boundary layer (W. O. Criminale, University of Washington);
- Analytical integration of the kinetic energy equation based on the perturbed velocity and vorticity field (*G. Staffilani*, *MIT*)
 - \Rightarrow Study of the intermediate and long term asymptotics.